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Chlorine Trifluoride Oxide. VII. The Difluorooxychloronium(V) Cation, ClF_2O^+ . Vibrational Spectrum and Force Constants

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The vibrational spectra have been recorded for the solid 1:1 adducts, $\text{ClF}_3\text{O} \cdot \text{BF}_3$, $\text{ClF}_3\text{O} \cdot \text{AsF}_5$, and $\text{ClF}_3\text{O} \cdot \text{SbF}_5$, and for $\text{ClF}_3\text{O} \cdot \text{BF}_3$ in HF solution. These spectra are entirely consistent with the ionic structures, $\text{ClF}_2\text{O}^+ \text{BF}_4^-$, $\text{ClF}_2\text{O}^+ \text{AsF}_6^-$, and $\text{ClF}_2\text{O}^+ \text{SbF}_6^-$, respectively. Six fundamental vibrations have been observed for ClF_2O^+ , consistent with symmetry C_s . The structure of ClF_2O^+ can be derived from a tetrahedron with the chlorine atom located at the center and with two fluorine atoms, one oxygen atom, and one free electron pair at the four corners. The vibrational spectrum and structure of ClF_2O^+ closely resemble those of isoelectronic SOF_2 . An assignment of the fundamental vibrations is proposed for ClF_2O^+ , and a modified valence force field has been calculated. A mainly covalent bond model is considered most likely for ClF_2O^+ .

Introduction

Chlorine trifluoride oxide possesses amphoteric character and forms adducts with BF_3 , AsF_5 , SbF_5 , and SiF_4 .¹ In this paper, we report the vibrational spectra and a structural study of some of these adducts.

Experimental Section

The preparation of the ClF_2O^+ containing salts, the apparatus, handling procedures, and the techniques used for recording the vibrational spectra have previously been described.¹⁻³

Results and Discussion

Vibrational Spectra.—Figures 1 and 2 show the infrared and Raman spectra of solid $\text{ClF}_2\text{O}^+ \text{BF}_4^-$ and $\text{ClF}_2\text{O}^+ \text{AsF}_6^-$, respectively. Figure 3 shows the Raman spectrum of $\text{ClF}_2\text{O}^+ \text{BF}_4^-$ in HF solution. The observed frequencies are listed in Table I and are compared with those reported for isoelectronic SOF_2 .⁴ The vibrational spectrum of the 1:1.2 $\text{ClF}_3\text{O} \cdot \text{SbF}_5$ adduct was also recorded. The infrared spectrum of a dry powder sample showed, in the range 4000–400 cm^{-1} , absorptions at the following wave numbers: 1339 sh, w, 1331 s, 1319 m, 740 s, 710 vs, 658 vs, 600 s, 559 mw, 510 ms, and 401 mw. The Raman spectrum of the solid showed bands at the following wave numbers (relative intensity): 1329 (4), 1316 (2), 741 (5), 686 (1), 670 (5), 637 (10), 590 (1), 554 (2), 504 (3), 397 (3), 377 (0+), and 282 br (4).

The adducts between ClF_3O and Lewis acids could be either ionic or fluorine bridged coordination complexes. The simplicity of the observed spectra, the shift of the $\text{Cl}=\text{O}$ stretching vibration to higher frequencies when compared to that in ClF_3O ,⁵ and the occurrence of all bands characteristic for the corresponding anion suggest that solid $\text{ClF}_3\text{O} \cdot \text{AsF}_5$ and $\text{ClF}_3\text{O} \cdot \text{BF}_3$ have the ionic structures $\text{ClF}_2\text{O}^+ \text{AsF}_6^-$ and $\text{ClF}_2\text{O}^+ \text{BF}_4^-$, respectively. The Raman spectrum of $\text{ClF}_2\text{O}^+ \text{BF}_4^-$ in HF solution closely resembles that of the solid, indicating the presence of discrete ions in both the solid state and HF solution. Table I shows the observed vibrational spectra, together with their assignment. The vibrational spectra of AsF_6^- and BF_4^- ,

respectively, are well known and hence will be discussed first.

An octahedral ion, such as AsF_6^- , has O_h symmetry. The six normal modes of vibration are classified as $A_{1g} + E_g + 2F_{1u} + F_{2g} + F_{2u}$. Of these, only the two F_{1u} modes will be infrared active, while only the A_{1g} , E_g , and F_{2g} modes will be Raman active, assuming that the selection rules are valid and that the octahedron is not distorted. The remaining F_{2u} mode is inactive in both the infrared and Raman spectrum. For $\text{ClF}_3\text{O} \cdot \text{AsF}_5$, all five active modes belonging to AsF_6^- were observed, the observed frequencies and intensities being similar to those of other AsF_6^- containing complexes, such as $\text{ClF}_2^+ \text{AsF}_6^-$,⁶ $\text{NF}_4^+ \text{AsF}_6^-$,⁷ $\text{NF}_2\text{O}^+ \text{AsF}_6^-$,⁸ $\text{IF}_6^+ \text{AsF}_6^-$,⁹ and $\text{ClO}_2^+ \text{AsF}_6^-$.¹⁰ Several deviations from the selection rules for O_h were observed. The $\nu_2(E_g)$ mode should be Raman active only but was also observed in the infrared spectrum. Furthermore, the two triply degenerate modes, ν_3 and ν_4 , in species F_{1u} show shoulders in the infrared spectrum. This breakdown of the selection rules has also been observed for the infrared spectra of most of the other AsF_6^- containing salts.⁶⁻¹⁰ It could be due to site symmetry lowering, slight distortion of the AsF_6^- octahedron, or weak fluorine bridges.

A tetrahedral ion, such as BF_4^- , has T_d symmetry. The four normal modes of vibration are classified as $A_1 + E + 2F_2$. Of these, all four modes are expected to be Raman active, whereas only the two F_2 modes should be infrared active. However, crystal-field effects or slight distortion of the BF_4^- tetrahedron can result in the A_1 mode also becoming infrared active. For $\text{ClF}_3\text{O} \cdot \text{BF}_3$, all four modes belonging to BF_4^- were observed and assigned by analogy with the known spectra of $\text{K}^+ \text{BF}_4^-$,¹¹⁻¹⁴ $\text{ClF}_2^+ \text{BF}_4^-$,⁶ $\text{ClO}_2^+ \text{BF}_4^-$,¹⁰ and $\text{FCl}_2^+ \text{BF}_4^-$.^{15,16}

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TABLE I
VIBRATIONAL SPECTRA OF ClF₂O⁺ SALTS COMPARED WITH THAT OF SOF₂

Observed frequencies, cm ⁻¹ , and relative intensities										
ClF ₂ O ⁺ BF ₄ ⁻			ClF ₂ O ⁺ AsF ₆ ⁻		SOF ₂ ^a			Assignment		
HF soln	Solid	Ir	Raman	Ir	Raman (gas)	Raman (liq)	Ir (gas)	OXF ₂ (C _s)	AsF ₆ ⁻ (O _h)	BF ₄ ⁻ (T _d)
1333 (3.5) } P 1323 sh	1333 (4) 1321 sh (2)	1334 s 1322 m 1295 m	1333 (2) 1320 (1)	1331 ms 1319 mw	1339 } (10) 1329	1308 (10)	1341 } s 1331			
	1020 (0+), br	990-1150 vs		820 sh						ν ₁ + ν ₄ (F ₂) ν ₂ (F ₂)
771 (2), P 741 (10), P 710 sh, dp	771 (3) 731 (10) 695 (5)	772 w 734 m 694 s	757 br (3) 696 (1)	750 br, s 695 vs 695 vs 675 sh, w 561 ms	808 (10) 747 (4)	804 (6) 716 (6)	808 s 747 vs			ν ₁ + ν ₃ (F _{2g}) ν ₁ (A ₁)
	530 sh (1) 519 (2)	532 w 521 sh								ν ₃ (F _{1u}) ν ₁ (A _{1g}) ν ₂ (E _g)
512 (2.5), P 402 (1.9)	513 (3) 404 (4)	514 s 405 m	511 (2) 406 (2)	509 ms 407 sh 388 s	530 (8) 390 (4)	528 (8) 399 (7)	530 w 393 w			ν ₃ (A') ν ₆ (A'') ν ₄ (F ₂)B ¹⁰ ν ₄ (F ₂)B ¹¹
383 (1)	384 (2)	383 m	378 sh (1) 371 (4)	378 sh	390 (4)	380 (5)	378 vw			ν ₄ (F _{1u}) ν ₄ (A') ν ₅ (F _{2g})
355 (0+)	356 (1)	356 mw								ν ₂ (E)

^a See ref 4.

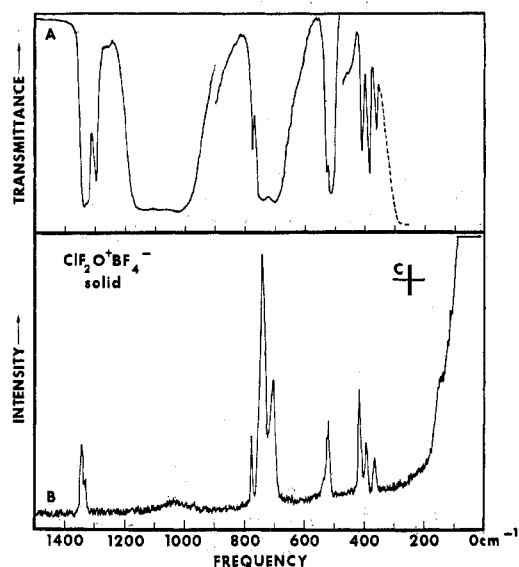


Figure 1.—Vibrational spectrum of solid ClF₂O⁺BF₄⁻: (A) infrared spectrum recorded as a AgCl disk; (B) Raman spectrum, exciting line 4880 Å. C indicates spectral slit width.

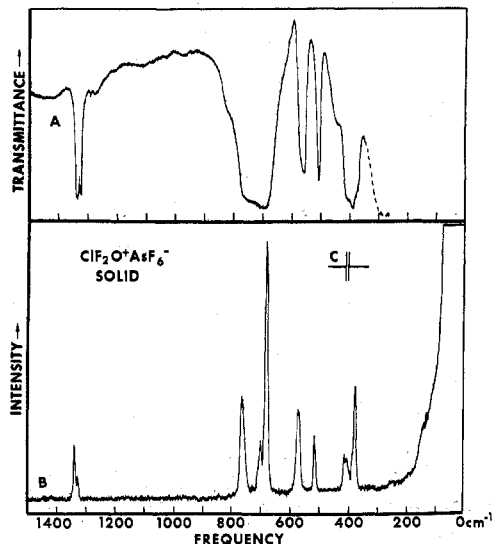


Figure 2.—Vibrational spectrum of solid ClF₂O⁺AsF₆⁻.

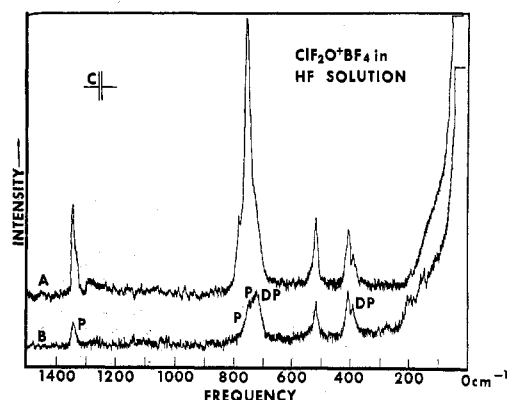


Figure 3.—Raman spectrum of ClF₂O⁺BF₄⁻ in HF solution: traces A and B, incident polarization perpendicular and parallel, respectively. C indicates spectral slit width.

The six remaining bands observed in the spectra of both ClF₂O⁺AsF₆⁻ and ClF₂O⁺BF₄⁻ complexes should be due to ClF₂O⁺. A four-atom ion of the type ZXY₂, such as ClF₂O⁺, could have C_{2v} or C_s symmetry. Distinction between symmetry C_{2v} and C_s should be possible by means of polarization measurements on the Raman bands. For symmetry C_{2v}, we would expect three polarized Raman bands, whereas for C_s, four of them should be polarized. The Raman spectrum of ClF₂O⁺BF₄⁻ in HF solution (Figure 3) shows that of the bands assigned to ClF₂O⁺, three bands (1333, 741, and 512 cm⁻¹) are clearly polarized and one (710 cm⁻¹) is depolarized. It is difficult to determine the polarization state of the remaining two lines at 404 and 384 cm⁻¹, respectively. However, polarization measurements⁴ for isoelectronic SOF₂, which has a Raman spectrum very similar to that of ClF₂O⁺, indicate that the less intense lower frequency mode is weakly polarized, and, hence, should be assigned to ν₄(A'). Intuitively, however, one might expect the ν₄(A') mode to result in a more intense Raman line than ν₆(A''). Unfortunately, Figure 3 does not allow a clear-cut distinction between the two possible assignments. Consequently, assignments for ClF₂O⁺ were made by complete analogy with SOF₂, which was shown¹⁷ to

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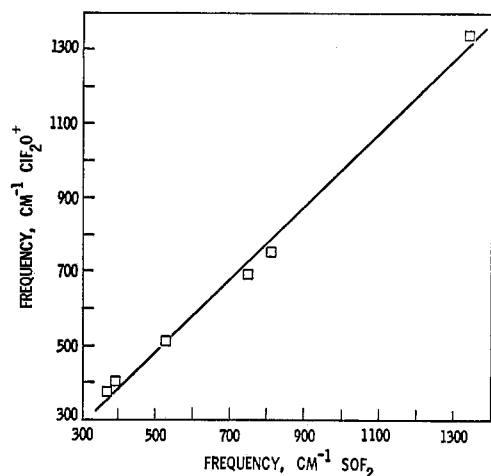


Figure 4.—Correlation between the vibrational frequencies of ClF_2O^+ and those of SOF_2 .

have symmetry C_s . A plot of the vibrational frequencies of ClF_2O^+ against those of SOF_2 (see Figure 4) shows an approximately linear correlation. This linearity, together with relative band intensity considerations and the results from the force constant calculations (see below), supports the assignments given in Table I, although a possible reversal of the assignments of $\nu_4(\text{A}')$ and $\nu_6(\text{A}'')$ cannot entirely be excluded.

The spectra of both ClF_2O^+ and SOF_2 species show a frequency splitting of $\nu_1(\text{A}')$. Whereas for SOF_2 the splitting is due to Fermi resonance⁴ between $\nu_1(\text{A}')$ and $\nu_2 + \nu_3(\text{A}')$, the splitting for ClF_2O^+ is due to the ³⁵Cl and ³⁷Cl isotopes. In ClF_2O^+ , the combination band $\nu_2 + \nu_3$ has too low a frequency to account for the splitting of the band at about 1330 cm^{-1} . Furthermore, the frequency splitting of 12.5 cm^{-1} , observed under high resolution conditions, agrees well with the value of 12.6 cm^{-1} calculated for ³⁵ ClF_2O^+ and ³⁷ ClF_2O^+ .

The vibrational spectrum obtained for the adduct $\text{ClF}_3\text{O} \cdot 1.22\text{SbF}_6^-$ clearly shows the bands characteristic for the ClF_2O^+ cation at $1331, 1319, 740, \sim 700, 510, 401,$ and 377 cm^{-1} . The remaining bands can be attributed to either SbF_6^- or polymeric anions such as $\text{Sb}_2\text{F}_{11}^-$.¹⁸⁻²⁰ Consequently, the formulation of the adduct as $\text{ClF}_2\text{O}^+\text{SbF}_6^-$ (containing some $\text{ClF}_2\text{O}^+\text{Sb}_2\text{F}_{11}^-$) appears to be appropriate. Similarly, the 1:1 adduct²¹ obtained in the reaction between ClF_3O and PtF_6^- contains the ClF_2O^+ cation and, hence, is best described as $\text{ClF}_2\text{O}^+\text{PtF}_6^-$. The infrared spectrum²¹ of $\text{ClF}_2\text{O}^+\text{PtF}_6^-$ shows, in the range $520\text{--}350\text{ cm}^{-1}$, no absorptions due to the anion and, hence, proves that the three ClF_2O^+ deformation modes occur at about $500, 400,$ and 380 cm^{-1} .

Force Constants.—Two sets of force constants were computed for ClF_2O^+ , assuming $\nu_6 > \nu_4$ (set I) and $\nu_4 > \nu_6$ (set II). Since SOF_2 and SeOF_2 are isoelectronic and have similar vibrational frequencies,⁴ their force constants were reevaluated by the same method for

comparison. The potential and kinetic energy metrics were computed with a machine method²² using the geometries shown in Table II. The ClO and ClF

TABLE II
BOND DISTANCES (Å) AND BOND ANGLES (DEG)
ASSUMED FOR ClF_2O^+ , SOF_2 , AND SeOF_2

	$\text{ClF}_2\text{O}^{\text{a}}$	SOF_2^{b}	SeOF_2^{c}
$R(\text{X}=\text{O})$	1.41	1.412	1.576
$r(\text{X}-\text{F})$	1.62	1.585	1.7255
$\beta(\text{O}=\text{X}-\text{F})$	108	106.82	104.82
$\alpha(\text{F}-\text{X}-\text{F})$	93	92.82	92.22

^a Estimated. ^b Reference 17. ^c I. C. Bowater, R. D. Brown, and F. R. Burden, *J. Mol. Spectrosc.*, **28**, 461 (1968).

bond lengths for ClF_2O^+ were estimated by comparison with similar molecules. The correlation noted by Robinson²³ between stretching frequencies and bond lengths predicts a value of 1.39 Å for the ClO bond length, which is only slightly shorter than that accepted here. The bond angles in ClF_2O^+ were estimated to be slightly larger than those in SOF_2 , assuming that the mutual repulsion between the ligands will increase with decreasing ionic radius of the central atom.

The force constants were computed by trial and error with the aid of a computer, requiring exact fit between the observed and computed frequencies. The results are given in Table III where the force constants

TABLE III
VIBRATIONAL FORCE CONSTANTS OF ClF_2O^+ , SOF_2 ,
AND SeOF_2^{a}

	ClF_2O^+		SOF_2	SeOF_2
	$\nu_6 > \nu_4$	$\nu_4 > \nu_6$		
f_R	11.21	11.20	10.84	7.87
f_r	3.44	3.44	4.03	3.57
f_β	1.72	1.65	1.73	1.28
f_α	1.59	1.78	1.52	1.04
$f_{\beta\beta}$	0.32	0.21	0.43	0.23
f_{rr}	0.25	0.39	0.42	0.37

^a Stretching constants in mdyn/Å and deformation constants in mdyn Å/radian^2 .

not shown were assumed to be zero. The values shown for the interaction constants, while not unique, were the simplest set that would give an exact fit. Since the force constants are underdetermined, a statistically meaningful uncertainty estimate cannot be made. Numerical experiments show that assuming different interaction constants, constrained to plausible values, resulted in valence force constants differing from those shown in Table III by less than 0.2. These uncertainties are estimates and are felt to be conservative. Of the two sets of force constants obtained for ClF_2O^+ , set I ($\nu_6 > \nu_4$) is more plausible since it results in f_β ($\angle \text{O}=\text{Cl}-\text{F}$) being larger than f_α ($\angle \text{F}-\text{Cl}-\text{F}$). The good agreement between the force constants of the three species, ClF_2O^+ , SOF_2 , and SeOF_2 , adds further credibility to our assignments for ClF_2O^+ . Table IV shows the potential energy distribution for set I of ClF_2O^+ . For set II, the results were similar, except for an appreciably larger contribution from $f_{\beta\beta}$ to ν_6 . The potential energy distribution for SOF_2 and SeOF_2 was quite similar to that of ClF_2O^+ . Only the

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TABLE IV
 POTENTIAL ENERGY DISTRIBUTION FOR ClF_2O^+

Assignment	Frequency, cm^{-1}	Potential energy distribution
A' ν_1	1333	$0.97f_R$
ν_2	731	$0.80f_r + 0.07f_\alpha + 0.06f_\beta + 0.06f_{rr}$
ν_3	513	$0.76f_\beta + 0.14f_{\beta\beta} + 0.05f_r$
ν_4	384	$0.91f_\alpha + 0.08f_r$
A'' ν_5	695	$0.96f_r + 0.14f_\beta - 0.07f_{rr}$
ν_6	404	$1.09f_\beta - 0.20f_{\beta\beta} + 0.11f_r$

more important terms are given in Table IV. These data show that the normal modes assigned to the observed vibrational frequencies are a good representation of the group vibrations described by the symmetry coordinates.

Force constants have previously been computed for both SOF_2 and SeOF_2 . Our values for SeOF_2 are very similar to those reported by Paetzold.²⁴ For SOF_2 , however, our results do not agree with those computed by Cotton and Horrocks.²⁵ The largest discrepancy was found for f_α ²⁶ for which Cotton and Horrocks²⁵ report an abnormally high value of 3.47 mdyn $\text{\AA}/\text{radian}^2$. This value appears much too high for this type of deformation vibration and cannot be explained by the slightly higher frequency value of 410 cm^{-1} assumed for ν_4 by Cotton and Horrocks.²⁵

The two constants of greatest interest are the two stretching force constants. Table V, showing a com-

 TABLE V
 Cl-O VALENCE FORCE CONSTANTS AS A FUNCTION OF THE OXIDATION NUMBER AND FORMAL ELECTRICAL CHARGE OF THE CENTRAL ATOM

Species	f_R , mdyn/ \AA	Oxidation state of the central atom	Formal electrical charge on the central atom
ClF_2O^+	11.21	+V	Positive
ClO_2F^a	9.41	+VII	
ClF_3O^b	9.37	+V	
$\text{ClF}_4\text{O}^{-c}$	9.13	+V	Negative
ClO_2F^d	9.07	+V	
ClO_2^+e	8.96	+V	Positive

^a W. Sawodny, A. Fadini, and K. Ballein, *Spectrochim. Acta*, **21**, 995 (1965). ^b Reference 5. ^c Reference 2. ^d D. F. Smith, G. M. Begun, and W. H. Fletcher, *Spectrochim. Acta*, **20**, 1763 (1964). ^e Reference 10.

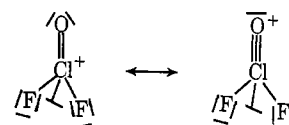
parison of the Cl—O valence force constants of similar molecules and ions, reveals that f_R in ClF_2O^+ is larger by about 1.8 mdyn/ \AA than the largest value known for any other Cl=O containing compound. Whereas the formal positive charge on the central atom in ClF_2O^+ will certainly contribute to the high f_R value,

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(26) It should be pointed out that the definition of the angles α and β in our work and that of ref 23 is reversed. Furthermore, the value of f_α in ref 23 was normalized for distance to allow a better comparison.

it cannot account for it solely. Contributions from the resonance structures below could explain the high



f_R value in ClF_2O^+ . The following observation seems noteworthy. In ClF_4O^- , the Cl—O valence force constant value² is about the same as those in similar compounds, but the Cl—F valence force constant decreases dramatically. In ClF_2O^+ , it is the ClO valence force constant which increases significantly, whereas the ClF valence force constant is within the usual range. This seems to indicate that the effective electronegativity of pentavalent chlorine is intermediate between those of oxygen and fluorine. Hence, in ClF_2O^+ the positive charge is partially located on the oxygen atom, whereas in ClF_4O^- the negative charge is mainly located on the fluorine ligands.

The relatively high values obtained for both the ClO and ClF valence force constants in ClF_2O^+ indicate that the bonding in ClF_2O^+ is best described by a mainly covalent bond model employing sp^3 hybridized orbitals of the chlorine atom for the two Cl—F bonds, the Cl—O σ bond, and the free electron pair. Bond models involving semiionic three-center, four electron p—p σ bonds^{27–30} are not possible owing to the geometry of ClF_2O^+ (absence of linear F—Cl—F groups) and would result in a considerably lower value for the Cl—F valence force constants.

In summary, the ClF_3O ·Lewis adducts investigated in this study are ionic and contain the ClF_2O^+ cation. The structure and bonding of this cation strongly resemble those of isoelectronic SOF_2 . The close correlation between the two isoelectronic couples, $\text{ClF}_2\text{O}^+ - \text{SOF}_2$ and $\text{ClF}_4\text{O}^- - \text{XeOF}_4$, respectively, demonstrates that the analogy between halogen fluorides and isoelectronic chalcogen or noble gas fluorides is likely to hold also for the oxyfluorides.

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